



(86) Date de dépôt PCT/PCT Filing Date: 1996/06/04
(87) Date publication PCT/PCT Publication Date: 1996/12/12
(45) Date de délivrance/Issue Date: 2007/08/07
(85) Entrée phase nationale/National Entry: 1997/10/27
(86) N° demande PCT/PCT Application No.: GB 1996/001324
(87) N° publication PCT/PCT Publication No.: 1996/039719
(30) Priorité/Priority: 1995/06/05 (GB9511336.1)

(51) Cl.Int./Int.Cl. *G02B 1/02* (2006.01),
G02B 5/08 (2006.01), *H01L 31/0232* (2006.01),
H01L 31/0352 (2006.01)

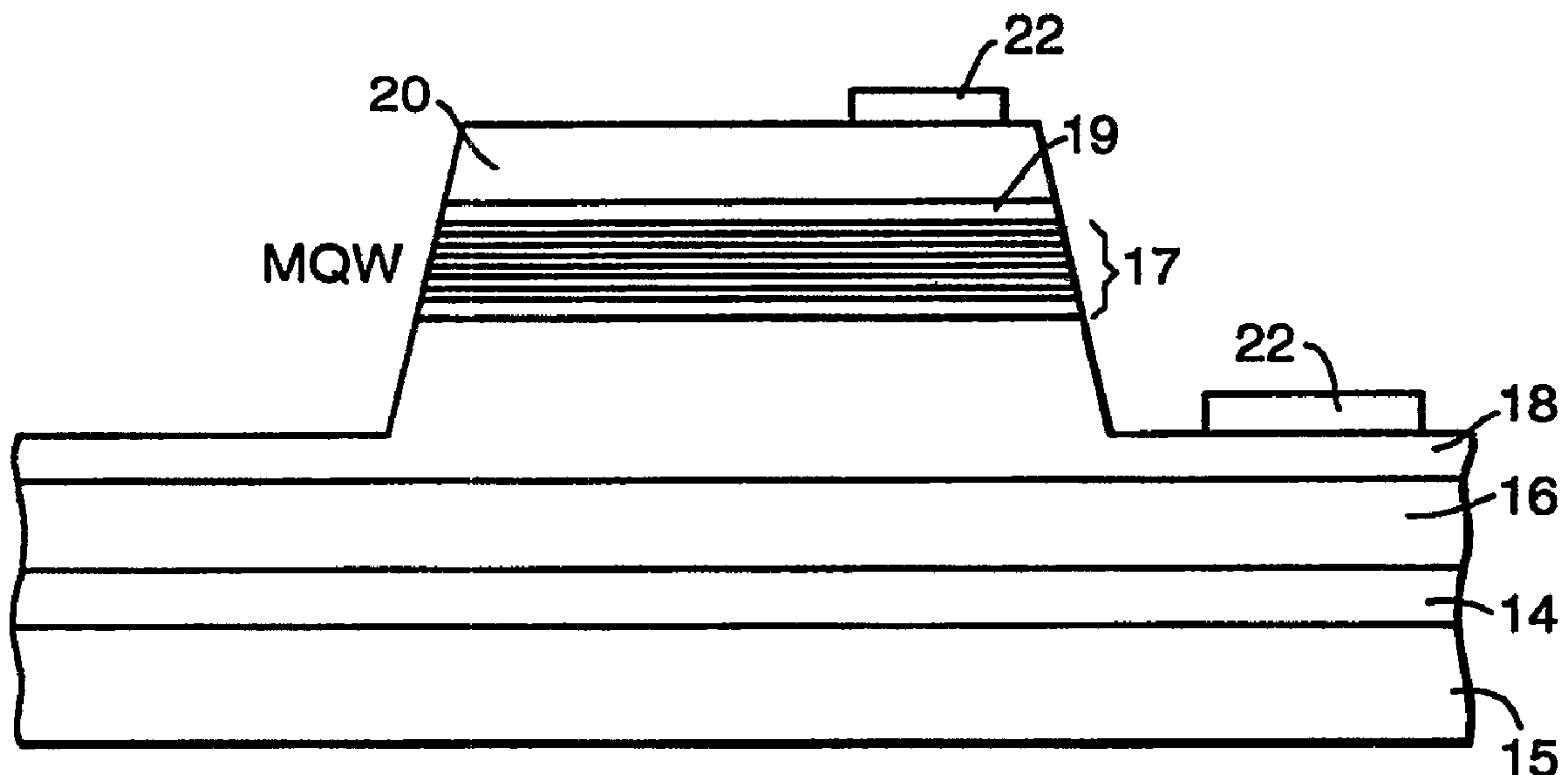
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(54) Titre : SUBSTRATS REFLECHISSANTS A SEMICONDUCTEURS

(54) Title: REFLECTING SEMICONDUCTOR SUBSTRATES



(57) Abrégé/Abstract:

A semiconductor substrate, suitable for epitaxial growth thereon, comprising a plurality of layers of material. The interfaces between layers act as reflectors of electromagnetic radiation. The reflectors may be used in, for example, resonant cavities in which may be located, for example, multi-quantum well detectors, the efficiency of said detectors being increased by virtue of the enhanced electric field associated with resonance in the cavity.



PCTWORLD INTELLECTUAL PROPERTY ORGANIZATION
International Bureau

INTERNATIONAL APPLICATION PUBLISHED UNDER THE PATENT COOPERATION TREATY (PCT)

(51) International Patent Classification ⁶ : H01L 31/0352, 31/0232	A1	(11) International Publication Number: WO 96/39719 (43) International Publication Date: 12 December 1996 (12.12.96)
(21) International Application Number: PCT/GB96/01324 (22) International Filing Date: 4 June 1996 (04.06.96) (30) Priority Data: 9511336.1 5 June 1995 (05.06.95) GB (71) Applicant (for all designated States except US): THE SECRETARY OF STATE FOR DEFENCE [GB/GB]; Defence Evaluation & Research Agency, DRA Farnborough, Hampshire GU14 6TD (GB). (72) Inventors; and (75) Inventors/Applicants (for US only): CARLINE, Roger, Timothy [GB/GB]; Defence Research Agency, St. Andrews Road, Malvern, Worcestershire WR14 3PS (GB). ROB-BINS, David, John [GB/GB]; Defence Research Agency, St. Andrews Road, Malvern, Worcestershire WR14 3PS (GB).		(81) Designated States: AM, AT, AU, BB, BG, BR, BY, CA, CH, CN, CZ, DE, DK, EE, ES, FI, GB, GE, HU, IL, IS, JP, KE, KG, KP, KR, KZ, LK, LR, LT, LU, LV, MD, MG, MN, MW, MX, NO, NZ, PL, PT, RO, RU, SD, SE, SG, SI, SK, TJ, TM, TT, UA, UG, US, UZ, VN, European patent (AT, BE, CH, DE, DK, ES, FI, FR, GB, GR, IE, IT, LU, MC, NL, PT, SE). Published <i>With international search report.</i> <i>With amended claims.</i>
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(57) Abstract A semiconductor substrate, suitable for epitaxial growth thereon, comprising a plurality of layers of material. The interfaces between layers act as reflectors of electromagnetic radiation. The reflectors may be used in, for example, resonant cavities in which may be located, for example, multi-quantum well detectors, the efficiency of said detectors being increased by virtue of the enhanced electric field associated with resonance in the cavity.		

REFLECTING SEMICONDUCTOR SUBSTRATES

The current invention relates to a method of realising reflectors of electromagnetic radiation in semiconductor substrates upon which high quality epitaxy can be grown. These substrates may be used to produce a resonant cavity for electromagnetic radiation: such cavities are of importance for applications such as laser etalons and infrared photon detectors.

Infrared detectors are of importance for both civilian and military applications. Where radiation is transmitted through the atmosphere the useful wavelength ranges are limited to $[3 - 5] \times 10^{-6}\text{m}$ and $[7.5 - 14] \times 10^{-6}\text{m}$ by atmospheric absorption. For military applications the long wavelength range may be further limited to about $[7.5 - 9] \times 10^{-6}\text{m}$ by filters designed to avoid dazzle by CO_2 lasers. Focal plane arrays (FPAs) with many individual detector elements are used for imaging in both these ranges. Several technologies are available for FPAs in the $[3 - 5] \times 10^{-6}\text{m}$ range eg CdHgTe (CMT) or InSb hybridised to Si readout circuits, or monolithic arrays. The technology is less well developed in the $[7.5 - 14] \times 10^{-6}\text{m}$ range. CMT is difficult to grow with the uniformity required for long wavelength infrared (LWIR) FPAs and requires cold bonding to the Si readout circuit which makes fabrication more difficult. Multi Quantum Well (MQW) detectors using GaAs/AlGaAs have progressed to 128×128 arrays using solder bump technology to hybridise to a Si readout chip but with unknown yield for larger arrays (B.F. Levine, J. Applied Physics 74 R1 (1993)). High performance FPAS are typically operated at temperatures close to that of liquid nitrogen. The thermal mismatch between detector and readout chips limits the size of hybrid arrays using compound semiconductor detectors on Si readout circuits. There is no successful Si monolithic LWIR FPA technology; IrSi, which has a long wavelength infrared response, has low quantum efficiency and therefore very low operating temperature.

At present LWIR technologies are being developed which are compatible with the Si readout circuits to avoid limitations inherent in the hybrid approach using compound semiconductor detectors. Of these the use of pseudomorphic heterostructures formed from silicon-germanium alloys epitaxially grown on Si substrates (SiGe/Si) is the most promising. Pseudomorphic SiGe/Si has the advantage of allowing a very high degree of uniformity over the Si substrate wafer which minimises fixed pattern noise in the FPA and favours high manufacturing yield.

LWIR photon detectors made from SiGe/Si MQWs are limited by low responsivity and by dark current thermal noise. Devices must be operated at lower temperatures to give acceptable signal/noise ratios. This restricts the utility and increases the cost of the imaging system. An increase in the quantum efficiency for absorption of incident radiation should improve the signal/noise ratio and so allow operation at higher temperature.

The quantum efficiency of detectors at shorter wavelengths ($< 2 \times 10^{-6} \text{m}$) has been enhanced using resonant cavities (see eg R Kuchibhotla, J Campbell, J C Bean, L Peticolas and R Hull, Appl. Phys. Lett 62 2215 (1993)). Resonant cavities produce localised regions of high electric field. Where absorption is dependent on the electric field the quantum efficiency can be enhanced by locating the absorbing region of a device in a resonant cavity.

To produce such a resonant cavity, the incident radiation, wavelength λ , is confined between two reflectors. Electric field enhancement in the cavity increases as the reflectances of the confining mirrors at wavelength λ increase. For a cavity of width L_c and refractive index n , light confined in the cavity resonates at wavelengths determined by optical thickness nL_c/λ , and the phase changes at the confining mirrors. Hence the cavity width, as well as the wavelength-dependence of the reflectors, must be chosen carefully to match λ . Provided the reflectors and absorption region are sufficiently broad-band, resonances at more than one wavelength may be used.

The reflectors used in resonant cavities are often dielectric or semiconductor stacks which consist of a large number of pairs of layers with refractive indices n_1 (higher than n for the cavity) and n_2 (lower than n for the cavity) and with individual layer thicknesses $\lambda/4n_i$ (where $i = 1$ or 2). An important property of such reflectors is their high transmission of incident light to allow penetration of the stacks and the cavity. Use of such Bragg reflectors for LWIR applications is limited, however, by difficulties in producing high quality pairs of thick layers to match the large wavelength. When one of the layers is strained, as for SiGe/Si, defects are introduced by strain relaxation in thick layers which reduce the efficiency of the Bragg reflector.

Semiconductor resonant structures are often grown epitaxially on single crystal substrates. A key requirement is then for a high efficiency reflector buried below the epitaxial active region of the device. Since epitaxial Bragg reflectors are unsuitable for LWIR applications, alternative types of buried reflector must be sought while maintaining the suitability of the substrate to produce high quality epitaxial structures. The reflector on top of the active region can simply be the semiconductor/air interface, or other layers grown epitaxially or deposited after epitaxy.

Single dielectric layers of SiO_2 , which have a refractive index of about 1.45 over a range of infrared wavelengths (whereas Si has a refractive index of about 3.45 in the infrared), have been suggested and demonstrated as buried reflectors in Si at shorter wavelengths (eg V P Kesan, P G May, F K LeGoues and S S Iyer, J. Cryst. Growth. 111 936 (1991); D K Nayak, N Usami, S Fukatsu and Y Shiraki, Appl. Phys. Lett. 64, 2373 (1994)). SIMOX wafers (separation by implantation of oxygen) were used as substrates. The high quality of Si overlying the oxide allows epitaxial overgrowth of Si and SiGe epilayers. The thickness of buried SiO_2 that can be produced in SIMOX wafers is, however, limited.

An alternative technology for the production of high quality buried oxides in Si is that of bond-and-etch silicon on insulator (BESOI), where limitations on thickness are less severe. This is important at longer wavelengths where the thicknesses of SiO₂ in the microns may be necessary to produce the required reflectances.

The weak absorption in SiO₂ at infrared wavelengths shorter than 7.5×10^{-6} m results in a wavelength dependent reflectance for a single layer which exhibits maxima and minima at wavelengths determined by the SiO₂ thickness. However, SiO₂ exhibits strong phonon absorption bands at about 9.2×10^{-6} m and about 21.5×10^{-6} m. It is a property of such absorption bands that, in a wavelength region on the short wavelength side of each peak, the refractive index shows a local minimum. As a result of this property, in these spectral regions the reflectance of a single thick SiO₂ layer buried in a Si ambient is enhanced compared with wavelength ranges of transparency removed from the absorption bands. In addition, in these spectral regions the reflectance is only weakly dependent on wavelength, and in the important wavelength region $[7.5-9] \times 10^{-6}$ m the reflectance is controlled by the oxide thickness, becoming almost independent of oxide thickness for layers $> 1.5 \times 10^{-6}$ m thick. This contrasts with the properties in the shorter wavelength range of dielectric transparency, where the reflectance is a strong function of both thickness and wavelength.

In principle, wafer bonding technology allows use of other buried dielectric layers having absorption bands, and enhanced reflectance, at different wavelengths. Similarly, a resonant cavity using a buried dielectric reflector and a different semiconductor material might be made by bonding that material to a Si wafer with the dielectric on its surface.

BESOI substrates have important advantages for making resonant cavity devices in the $[7.5-9] \times 10^{-6}$ m wavelength region. First, a thick ($> 1.5 \times 10^{-6}$ m) oxide layer can be used as the back reflector in the device, giving high reflectance. Secondly, the reflectance of a thick layer is not very sensitive to variations in thickness. Thirdly, the reflectance is reasonably constant over the wavelength region so that strong resonance can be observed even if the cavity width differs slightly from the design value, L_c . As a result, the device performance is more tolerant to variations in both oxide thickness and cavity width than would be the case at shorter wavelengths.

Buried layers of SiO_2 may also be used as front reflectors with light incident on the Si/ SiO_2 layer from the substrate side of the cavity. In this case the useful wavelength region of enhanced reflectance near absorption bands will be reduced by absorption of incident light in the SiO_2 layer, for which the optimum thickness will be less than for oxide used as a back reflector. Anti-reflection coatings may be used to increase the light entering the substrate.

S Fukatsu, D K Nayak and Y Shiraki, in Applied Physics Letters, 65, 3039 (1994), report the use of single layers of SiO_2 to provide reflectors for resonant cavities at wavelengths below 2×10^{-6} m. However, no reference is made to the spectral dependence of the refractive index, n , and extinction coefficient, k . Figure 2 of this reference shows the reflectance of SiGe/Si epilayers over a single SiO_2 layer buried in a Si ambient.

The wavelength dependence of the reflectance exhibits an oscillatory behaviour corresponding to two resonances, one in the overlying SiGe/Si layers and one in the SiO_2 reflector itself. The latter resonance occurs due to the low value of extinction coefficient in the SiO_2 layer which results in light being reflected from the lower SiO_2 /Si interface to give interference in the SiO_2 layer and so a strong wavelength dependence of the layer's reflectance.

Fukatsu et al explain their reflectance results using Si and SiO₂ refractive indices of $n_{\text{Si}}=3.45$ and $n_{\text{SiO}_2}=1.45$ respectively and by assuming low absorption (i.e. k is very small). Extending their model to longer wavelengths predicts a continuing oscillatory behaviour of the reflectance. In the important wavelength region $[7.5-9] \times 10^{-6}\text{m}$ their model fails to predict the enhanced magnitude and wavelength-and thickness insensitivities of the reflectance which occur near regions of high k values and which are exploited in the present invention.

According to this invention a reflector for electromagnetic radiation within a wavelength range comprises

a semiconductor substrate of first material,

a reflecting layer of second material, said second material having a refractive index whose variation with wavelength exhibits at least one local minimum within the range and

a layer of third semiconductor material which may be the same or different from the first material;

said layers being arranged so that the reflecting layer is sandwiched between the other two layers.

In a preferred embodiment the second material has an extinction coefficient whose variation with wavelength exhibits at least one local maximum associated with the local minimum in refractive index.

In a further preferred embodiment, the local minima and local maxima arise from a phonon resonance.

In a further preferred embodiment the wavelength range is within $[7.5 \text{ to } 14] \times 10^{-6}\text{m}$.

In a further preferred embodiment the wavelength range is within $[7.5 \text{ to } 9] \times 10^{-6} \text{m}$.

In a further preferred embodiment the first material is silicon and the second material is silicon dioxide.

In a further preferred embodiment the reflector is incorporated in a resonant cavity.

In a further preferred embodiment, the resonant cavity is incorporated in a detector for electromagnetic radiation.

In a further preferred embodiment the detector includes a multi-quantum well region able to absorb incident radiation, said radiation having an electric field vector with a component which is parallel with the layers and having a wavelength which approximates to the resonant wavelength of the cavity.

The invention will now be described, by example only, with reference to the following figures in which:-

Figure 1 shows the variation of refractive index and extinction coefficient with wavelength for SiO_2 ;

Figure 2 shows the predicted reflectance of SiO_2 , in Si ambient, when assumed to be an idealised dielectric with constant refractive index;

Figure 3 shows the actual variation of reflectance of SiO_2 , in Si ambient, with wavelength for a range of thicknesses as calculated from dielectric function spectra;

Figures 4 and 5 each show the depth dependence of electric field in a structure containing a Si cavity;

Figure 6a shows a schematic representation of a resonant cavity detector of the current invention (not to scale) and figure 6b shows the depth dependence of electric field of the device;

Figure 7 shows a second representation of a resonant cavity detector of the current invention (not to scale);

Figure 8 shows reflectance spectra for IR radiation normally incident on the top surface of a resonant cavity and a control device;

Figure 9a shows absolute responsivity spectra for the resonant cavity detector and control illuminated from the top surface, and reflectance for the resonant device, and Figure 9b shows absolute responsivity spectra and reflectance for a resonant cavity detector illuminated from the substrate side and

Figure 10 shows the absolute responsivity and reflectance spectra for a second resonant cavity detector designed to provide a peak in responsivity in the useful wavelength range of $[8 - 9] \times 10^{-6}\text{m}$.

Referring to figure 1, the variation of index of refraction n and extinction coefficient k of SiO_2 with wavelength are seen. Localised minima 1 in the refractive index are seen at wavelengths of about $9 \times 10^{-6}\text{m}$ and about $20 \times 10^{-6}\text{m}$. Localised maxima 2 in the extinction coefficient are seen at wavelengths of about $9.2 \times 10^{-6}\text{m}$ and $21.5 \times 10^{-6}\text{m}$. These data are taken from the Handbook of Optical Constants of Solids, E D Palik (ed.) (Academic Press) (1985).

Figure 2 shows the calculated reflectance of SiO_2 , when assumed to be an idealised dielectric (without localised absorption bands), as a function of wavelength for a range of thicknesses d_{ox} from $5 \times 10^{-7}\text{m}$ to $2 \times 10^{-6}\text{m}$. The refractive index is taken as constant and assigned the value characteristic of SiO_2 in short wavelength regions of transparency.

For wavelengths within $[7.5 - 9] \times 10^{-6}\text{m}$ the reflectance varies strongly with wavelength and thickness and shows a maximum value of only about 0.5.

Referring to figure 3, the actual variation of reflectance of SiO_2 with wavelength for a range of thicknesses d_{ox} from $5 \times 10^{-7}\text{m}$ to $2 \times 10^{-6}\text{m}$ is shown. For $d_{\text{ox}} > 1.5 \times 10^{-6}\text{m}$, the reflectance of SiO_2 is almost independent of the oxide thickness for wavelengths between $[7.5 - 9] \times 10^{-6}\text{m}$, with a maximum value of about 0.76 in this wavelength region. These data are calculated using J A Woollam Company, V.A.S.E. Integrated Data Acquisition and Analysis Software for Variable Angle Spectroscopic Ellipsometry, Version 1.3, with reference dielectric function spectra obtained from the Handbook of Optical Constants of Solids referenced above.

Figure 4 shows the calculated depth dependence of electric field for radiation of wavelength $8 \times 10^{-6}\text{m}$ in a $2.4 \times 10^{-6}\text{m}$ thick Si cavity. Regions 3, 4, 5 and 6 of the graph correspond to layers of air, silicon, SiO_2 and silicon respectively. Thus the front (light incident) reflector is a Si/air interface and the back reflector is a SiO_2 ($2 \times 10^{-6}\text{m}$)/Si interface.

The horizontal axis of the graph represents optical thickness, where that of each layer in the structure is given by $n_i d_i / \lambda$, where n_i and d_i are the refractive index and physical thickness of layer i , and λ is the resonant wavelength.

Figure 5 shows the calculated depth dependence of electric field for radiation of wavelength $8 \times 10^{-6}\text{m}$ in a $1.83 \times 10^{-6}\text{m}$ thick Si cavity. Regions 7, 8, 9, 10 and 11 of the graph represent layers of silicon, SiO_2 , silicon, metal and air respectively. Thus the front reflector is a SiO_2 ($0.8 \times 10^{-6}\text{m}$)/Si interface and the back reflector is a Si/Al interface.

Calculations to obtain figures 4 and 5 were made using commercial software written by H. A. Macleod (Thin Film Design Software for Apple Macintosh Version 2.1 (1983-1990)) and the reference dielectric function spectra previously referred to.

The use of buried SiO₂ reflectors for resonant cavity applications will now be illustrated by a description of the production of a resonant cavity infrared detector for radiation in the range $[7.5 - 9] \times 10^{-6}$ m. The detector uses multiple p-SiGe/Si quantum wells to absorb the incident radiation and so convert the incident radiative energy into an electrical signal. The purpose of the resonant cavity is to enhance the coupling of the incident light into the absorbing MQW within a desired wavelength range. This is done by locating the MQW in a region of enhanced electric field produced by the resonant cavity. The wavelength of the peak in the MQW absorption, the total cavity width and the reflectors must be chosen carefully to match the incident radiation.

In the present example this was achieved using detailed theoretical modelling to provide the device design and was realised using epitaxial growth of the cavity, including the SiGe/Si MQW and doped electrical contacts, on a commercially purchased BESOI substrate wafer obtained from Shin-Etsu Handotai Co., Ltd.

Simple device structures were then fabricated and their electrical response to spectral radiation measured using the photoconductivity technique. These results were then compared with the optical properties of the resonant cavity measured by spectral reflectance in order to show an enhancement in device performance at the resonant wavelength of the cavity when compared with a non-resonant control device.

The substrate comprised a 2×10^{-6} m SiO₂ layer buried under a layer of lightly-doped Si (about 10^{-6} m). The cavity width was determined by simulating the electric field profiles in the proposed resonant cavity in a similar way to that used to give figures 4 and 5, using the same software package and reference dielectric function spectra previously referred to, and allowing for the effects of doping (in an approximate manner) in both the SiGe and Si layers (D K Schroder, R N Thomas and J C Swartz, IEEE Transactions on Electron Devices, Vol ED-25 No 2 (February 1978) p254).

The design of the MQW was chosen from measurements made on an earlier non-resonant device which gave a photoconductive response to radiation above $5 \times 10^{-6}\text{m}$ which cut off between $[10 - 11] \times 10^{-6}\text{m}$ and with a peak in its spectral response around $8 \times 10^{-6}\text{m}$. In one resonant device the top Si/air interface formed the front reflector and the buried oxide formed the back reflector. In this case, finding the total cavity width for resonance at $8 \times 10^{-6}\text{m}$ (which includes the MQW, contact layers and Si spacers) was assumed to be equivalent to determining the correct thickness of an antireflection coating on a metal, where the "metal" is SiO_2 and the antireflection coating is the cavity above it (see S L Daryanani and G W Taylor, Optical and Quantum Electronics 25 (1993), p123).

This can be done by using the software to minimise the reflectance of monochromatic incident radiation by altering the thickness of the resonant cavity. The thickness chosen was the thinnest cavity in which the contact layers could be accommodated. The active region was a 20-period MQW comprising Si barrier layers about $50 \times 10^{-9}\text{m}$ thick and $\text{Si}_{0.87}\text{Ge}_{0.13}$ quantum wells about $8 \times 10^{-9}\text{m}$ thick containing acceptors at a concentration of about $1 \times 10^{16}\text{m}^{-2}$.

Referring to figure 6a, a multi-quantum well IR resonant cavity detector featuring a reflector of the current invention comprises a BESOI region 12 and a multi-quantum well region 13. Region 12, in turn, comprises a reflector layer 14 of SiO_2 sandwiched between a p^- doped Si substrate 15 and a Si layer 16 which is suitable for further epitaxial growth.

Region 13 comprises multiple layers of SiGe/Si 17 sandwiched between a p^+ doped Si layer 18 and a layer of Si 19 suitable for further epitaxial growth. This is topped with a layer 20 of p^+ doped silicon suitable for the attachment of an electrical contact (not shown). Optimum resonant enhancement is achieved with radiation of suitable wavelength normally incident on the top surface (Si/Air) surface, but absorption will be enhanced for radiation at non-normal incidence which is refracted close to the normal inside the cavity.

Referring to figure 6b, the electric field profile of the device of figure 6a, as obtained using the commercial software previously referred to, is shown.

As in figures 4 and 5, the horizontal axis represents optical thickness.

The maximum in the electric field is attenuated by approximately 25% by absorption in the contact layers and has a peak magnitude between 30Vm^{-1} and 35Vm^{-1} . The peak electric field in an equivalent Si substrate is about 12.5Vm^{-1} , indicating an enhancement in the electric field intensity of between 5 and 8 times, as given by the square of the ratio of the electric fields in the resonant and non-resonant cases. In this example peak 21 is coincident with the Multi Quantum Well region 17 shown in figure 6a.

Referring to figure 7, like numerals are used to represent layers which correspond with those shown in figure 6a. Aluminium electrical contacts 22 are fixed to p^+ doped layers 18 and 20.

A second resonant device was made to detect radiation normally incident from the substrate side, using the same epitaxial structure. In this case the $2 \times 10^{-6}\text{m}$ buried oxide formed the front reflector for the cavity, and Al was deposited on the top Si surface of the epitaxial layers to form the back reflector. The surface of the substrate exposed to the incident radiation was polished to improve transmission, but no anti-reflection coating was applied.

An equivalent active MQW and contact structure was epitaxially grown on a p^+ -Si substrate and used to make a non-resonant control device.

The active device layers were grown by Low Pressure Vapour Phase Epitaxy (LPVPE) in a stainless steel reactor with ultra-highvacuum background. The source gases used were silane and germane for growth of the Si and SiGe epilayers with diborane included for the p-type doping. These were mixed with hydrogen, typically at pressures of 13 - 130Pa. Surface oxide was desorbed at 900 - 920°C and a $5 \times 10^{-8}\text{m}$ Si buffer was deposited at 750°C. Growth of the device layers was nominally at 610°C.

Referring to figure 8 the presence of optical resonance in the device cavity was established by observing minima in the spectral reflectance of the device. Reflectance spectra on as-grown resonant cavity and control devices were measured using a Perkin-Elmer 983G infrared spectrophotometer with gold as 100% standard. Light normally incident on the top surface was used and spectra were taken from the centre 23, centre + 0.01m 24 and centre + 0.02m 25 of the wafer. Wavelengths where resonance occurs were identified as those regions 26 where the reflectance of the resonant cavity was below that of the control 27. As the location of the measurement is moved further from the centre of the resonant device wafer the resonances shift to lower wavelength.

This is due to a change in the cavity width across the substrate caused by a radial variation in the wafer temperature during growth.

Both the presence of multiple resonances and their tunability using the width of the resonant cavity are demonstrated in figure 8.

The device design was intended to give a cavity resonance centred at 8×10^{-6} m near the middle of the substrate. In figure 8 this is seen to have shifted to 7.25×10^{-6} m 26b. This is due to the presence of the buried SiO₂ layer in the substrate which alters the emissivity of the BESOI wafer compared with a bulk Si substrate. The emissivity is used to measure the wafer temperature during epilayer deposition and so the resonant cavity device was grown at a temperature slightly below the intended 610°C. As a result the cavity is thinner than intended which gives a resonance at shorter wavelength.

The thickness of the cavity near the centre of the substrate was measured by modelling the reflectance spectra using the commercial software previously described with the layer refractive indices for undoped Si and SiO₂ used to give the device design. The thickness obtained was about 3.2×10^{-6} m, and was $.37 \times 10^{-6}$ m less than the design value. Using this revised cavity thickness the electric field profile in the device structure was modelled at each of the wavelengths where reflectance minima were observed. All gave rise to an enhanced electric field in the cavity due to resonance, the greatest enhancement being at 7.25×10^{-6} m.

In order to reposition the principal resonance for top surface illumination within the desired wavelength range $[7.5-9] \times 10^{-6}$ m, it is simply necessary to adjust the epitaxial growth conditions to produce a cavity of thickness 3.57×10^{-6} m. Alternatively, the resonance phase conditions within the thinner cavity can be changed by illuminating from the substrate side and using an Al back reflector deposited on the top Si surface.

The corresponding cavity resonance close to the middle of the wafer is then centred near 7.8×10^{-6} m.

In order to demonstrate the operation of resonant cavity detectors based on these principles, simple photoconductive devices were fabricated from both the resonant cavity and control samples. The devices were mesa structures with the electrical contacts made using evaporated Al. The following steps were used to produce the working cavity device with oxide back reflector and the control device for top surface illumination:-

- i) A piece of each sample, obtained by cleaving the wafers, had its surface oxide removed using dilute HF solution. This step is repeated before mesa etching and before Al evaporation. The sample may be immersed in methanol to inhibit surface reoxidation prior to subsequent process steps.
- ii) Wax dots were then applied to the surface to form etch masks of about 2mm diameter. These define the mesa area.

- iii) Approximately 3×10^{-6} m high mesas were formed using an aqueous etch comprising 30ml 0.825M Cr_2O_3 solution and 20ml 40% HF.
- iv) The wax masks were removed and Al evaporated on to the sample surface.
- v) Two wax dots were then applied to the Al to define two electrical contacts. On each device, one contact is on top of the mesa surface leaving the majority of the mesa area clear to allow the incident light to enter the cavity. The second is on the etched surface to make electrical contact below the MQW active region.
- vi) Al etch was then used to remove the unwanted Al. The wax was then removed.
- vii) The samples were then heated at 385°C for 5 minutes in a nitrogen atmosphere to give ohmic contacts by alloying in the Al.
- viii) The samples were then mounted on individual packages and electrical contact was made by ball bonding the Al contacts to the appropriate part of the package.

Referring to figure 9a, photoconductivity measurements of the resonant cavity and control devices with top surface illumination were made at a temperature of 20K. Relative responsivity spectra were obtained by dividing the spectral response of each sample to monochromatic infrared radiation by the response to the same radiation of a pyroelectric detector which is assumed to have a constant response over the wavelength range of interest.

The monochromatic radiation was provided by a global source passed through a grating monochromator and an appropriate long pass filter. The radiation was chopped. The electrical signal from the sample was fed, via a preamplifier into a lock-in amplifier with the chopper frequency of the incident radiation used as reference frequency. The signal from the pyroelectric was fed directly into the lock in amplifier. The spectral dependence of the lock-in output was recorded as the raw photoconductivity spectrum.

Absolute responsivity spectra were obtained from the relative responsivity using a photocurrent measurement with a 500K black body oven. Absolute responsivity spectra for both the resonant cavity device 28 and the control 29 are shown in figure 9a where they are compared with the reflectance 30 from the resonant cavity sample. This figure clearly demonstrates an infrared detector with resonantly enhanced operation. A peaked structure and a strong enhancement in responsivity are seen for the resonant cavity device when compared with the control sample.

The $7.25 \times 10^{-6}\text{m}$ peak 31 in responsivity of the resonant cavity device can be seen to be caused by cavity resonance of incident radiation of this wavelength by comparison with the reflectance measured on the same structure. A local minimum 32 in reflectance is seen at that wavelength. The responsivity peak 31 due to resonance is stronger than peaks 33, 34, and 35 at other wavelengths due to the enhanced reflectance of the SiO_2 reflector associated with the local minimum in its refractive index.

The following steps were used to produce a working cavity device with Al back reflector and buried oxide front reflector for substrate-side illumination:-

- i) A piece of the resonant cavity sample on the BESOI substrate had the rough bottom surface of the substrate polished to improve transmission. This polished surface was protected by wax.

- ii) Oxide was removed from the top surface using dilute HF solution. This step is repeated before mesa etching and before Al evaporation. The sample may be immersed in methanol to inhibit surface reoxidation prior to subsequent process steps.
- iii) Al was evaporated on the top surface and wax dots were then applied to the surface to form etch masks of about 2mm diameter. These defined the mesa area and the top contact covering the whole of that area.
- iv) The exposed Al was etched away and approximately 3×10^{-6} m high mesas were formed using an aqueous etch comprising 30ml 0.825M Cr_2O_3 solution and 20ml 40% HF.
- v) A second Al layer was evaporated onto the surface and a wax dot applied to the etched area to define the bottom contact.

Steps vi), vii) and viii) described above were then followed, with the addition that a hole was drilled in the mounting package to allow illumination from the polished substrate side.

Referring to figure 9b, absolute responsivity measurements 36 made as described for figure 9a at 20K are compared with reflectance spectra 37 for substrate-side illumination of the resonant cavity. The reflectance 37 is greater than in figure 9a largely due to the highly reflective Al back layer. There is a minimum in reflectance 38 at 7.8×10^{-6} m which shows that the combination of SiO_2 front reflector and Al back reflector has moved the cavity resonance to longer wavelength compared with the minimum 32 in figure 9a.

The photoconductive responsivity 36 in figure 9b associated with the cavity resonance shows a peak 39 at 7.5×10^{-6} m, within the desired wavelength range. The absolute magnitude of this peak responsivity is greater than for the peak 31 in the top illuminated cavity device in figure 9a, due to the greater enhancement of electric field intensity for the SiO_2/Si and Si/Al reflector pair compared with the SiO_2/Si and Si/air reflector pair.

This effect is evident in the higher calculated peak electric field within the cavity in figure 5 compared with figure 4.

The responsivity peak 39 in figure 9b does not occur at the same wavelength as the minimum 38 in reflectivity 37, due to attenuation of the incident light by absorption in the 2×10^{-6} m thick front SiO₂ reflector. Further enhancement of peak responsivity should be achievable by optimisation of the dielectric front reflector in this configuration.

Figure 10 shows photoconductivity measurements for another device fabricated in a similar fashion to that from which figure 9a was derived. This device had a cavity thickness of approximately 4.1×10^{-6} m and was illuminated at the top surface. The absolute responsivity spectrum 40 of the device is shown along with both measured reflectivity 41 and simulated reflectivity 42 (by similar methods to that described previously). Measurements were made at 2V bias.

Figure 10 shows the device to have a peak 43 in responsivity in the useful wavelength range of $[8 - 9] \times 10^{-6}$ m. Peak 43 is coincident with a minimum 44 in reflectivity.

1. A reflector for electromagnetic radiation within a wavelength range comprising:

a semiconductor substrate of first material,

a reflecting layer of second material, said second material having a refractive index whose variation with wavelength exhibits at least one local minimum within the range and an extinction coefficient whose variation with wavelength exhibits at least one local maximum associated with the local minimum in refractive index, and

a layer of third semiconductor material which may be the same or different from the first material;

said layers being arranged so that the reflecting layer is sandwiched between the other two layers.

2. The reflector of claim 1 where the local minima and local maxima arise from a phonon resonance.

3. The reflector of claim 2 where the wavelength range is within $[7.5 \text{ to } 14] \times 10^{-6} \text{ m}$.

4. The reflector of claim 3 where the wavelength range is within $[7.5 \text{ to } 9] \times 10^{-6} \text{ m}$.

5. The reflector of claim 4 where the first material is silicon and the second material is silicon dioxide.

6. A resonant cavity incorporating reflector according to any of claims 1 - 5.

7. A detector of electromagnetic radiation incorporating a resonant cavity according to claim 6.

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8. A detector according to claim 7 and further including a multiquantum well region able to absorb incident radiation, said radiation having an electric field vector with a component which is parallel with the layers and having a wavelength which approximates to the resonant wavelength of the cavity.

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Fig.1.

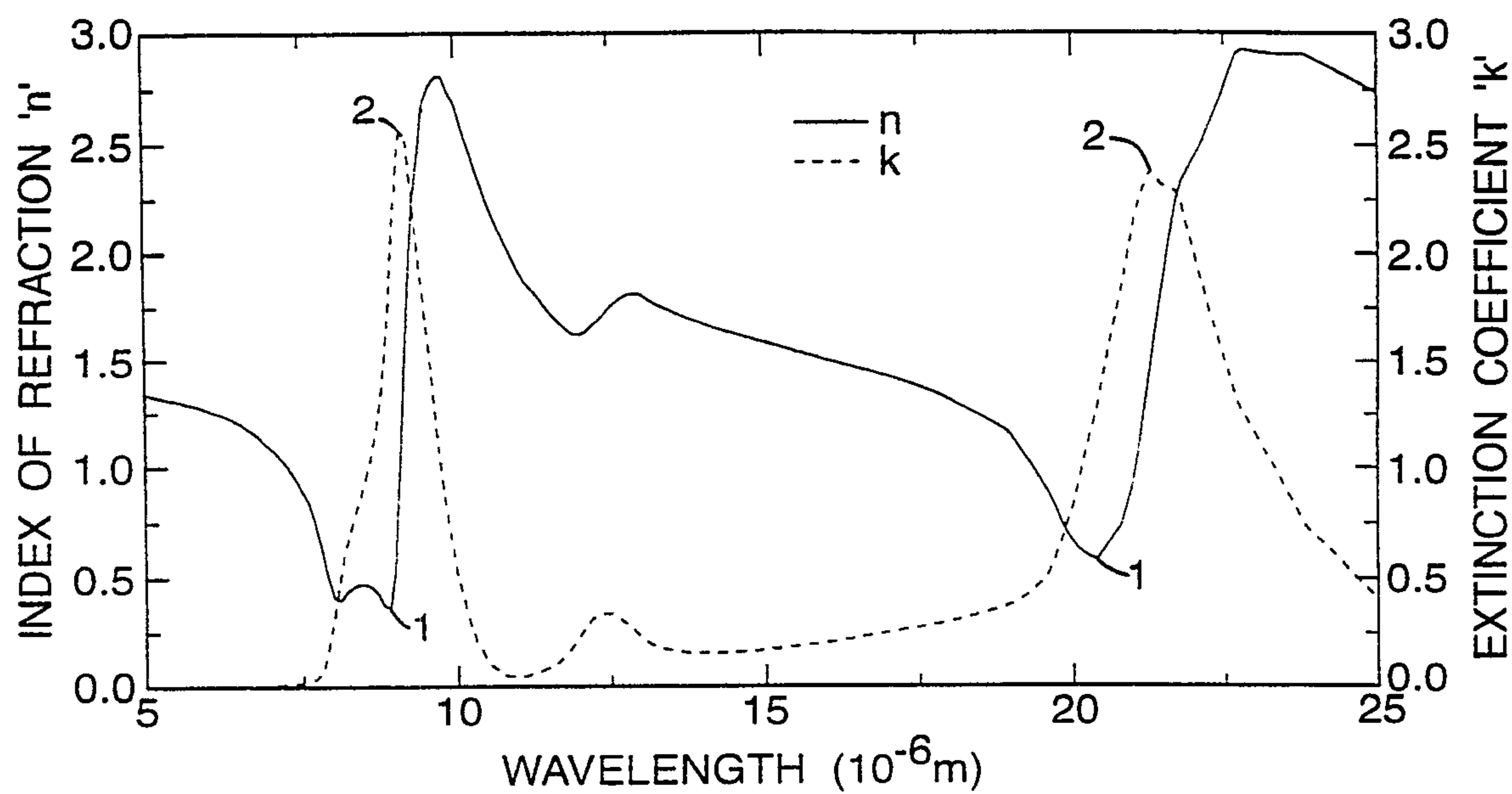
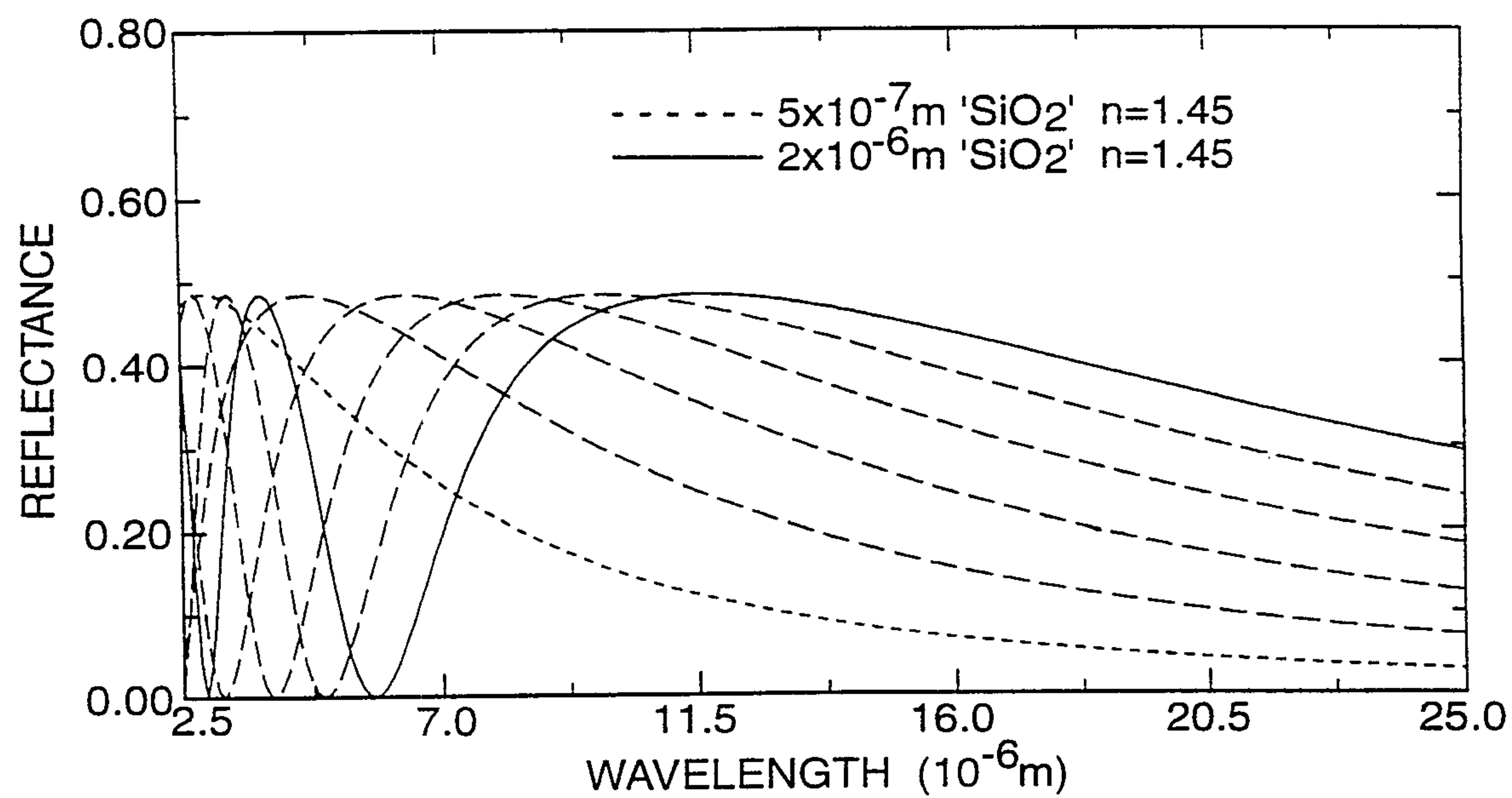


Fig.2.



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Fig.3.

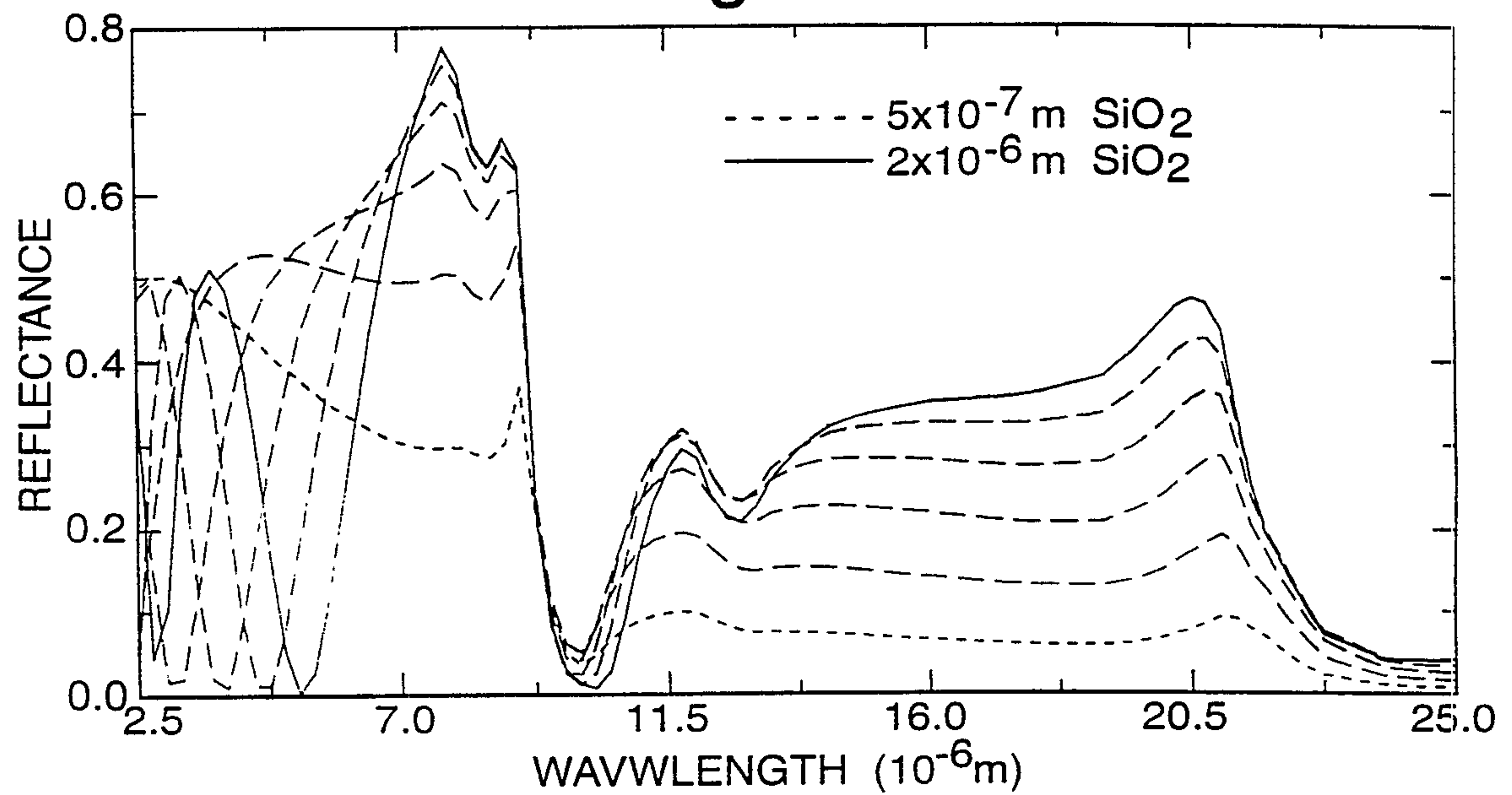
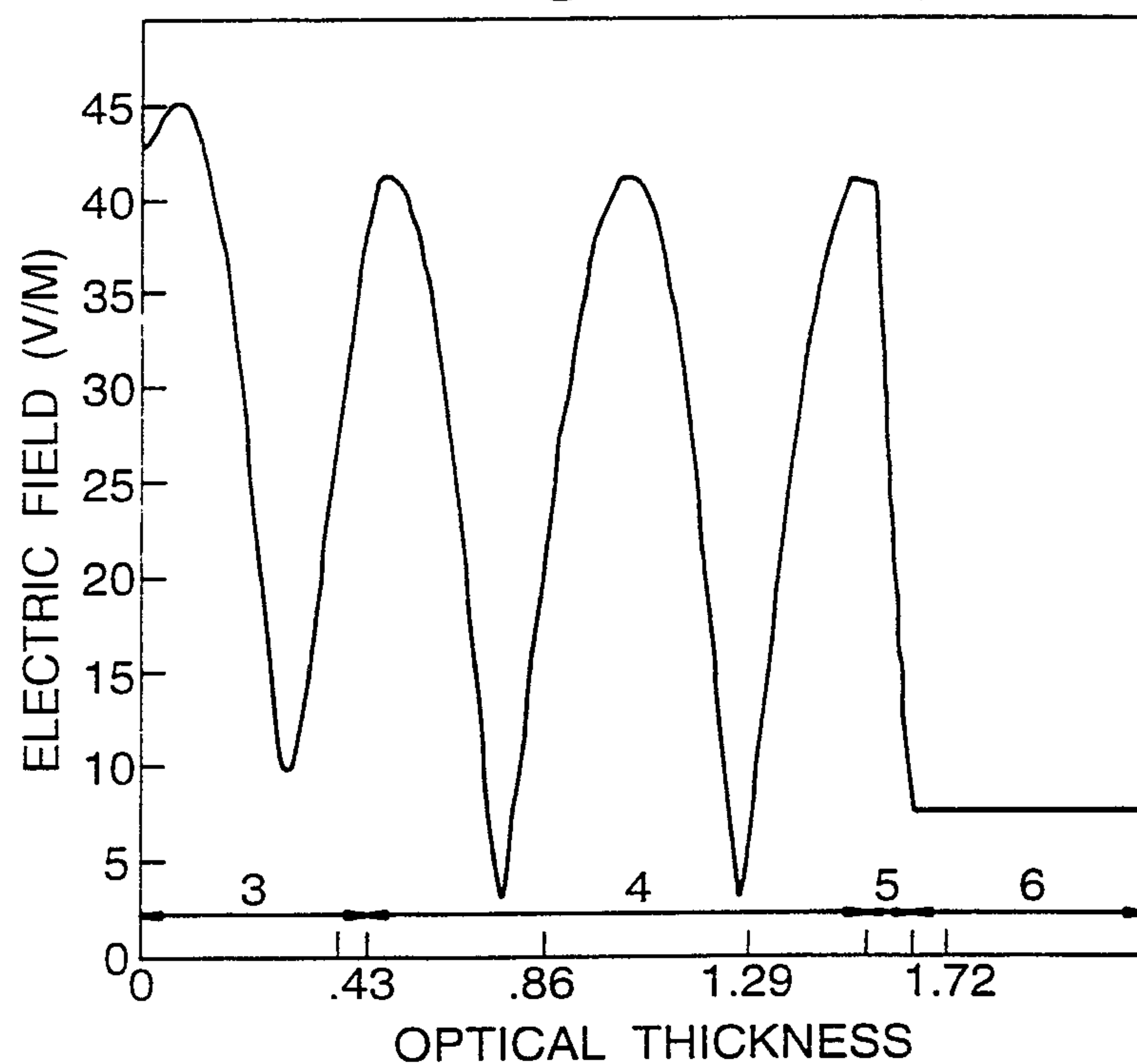


Fig.4.

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Fig.5.

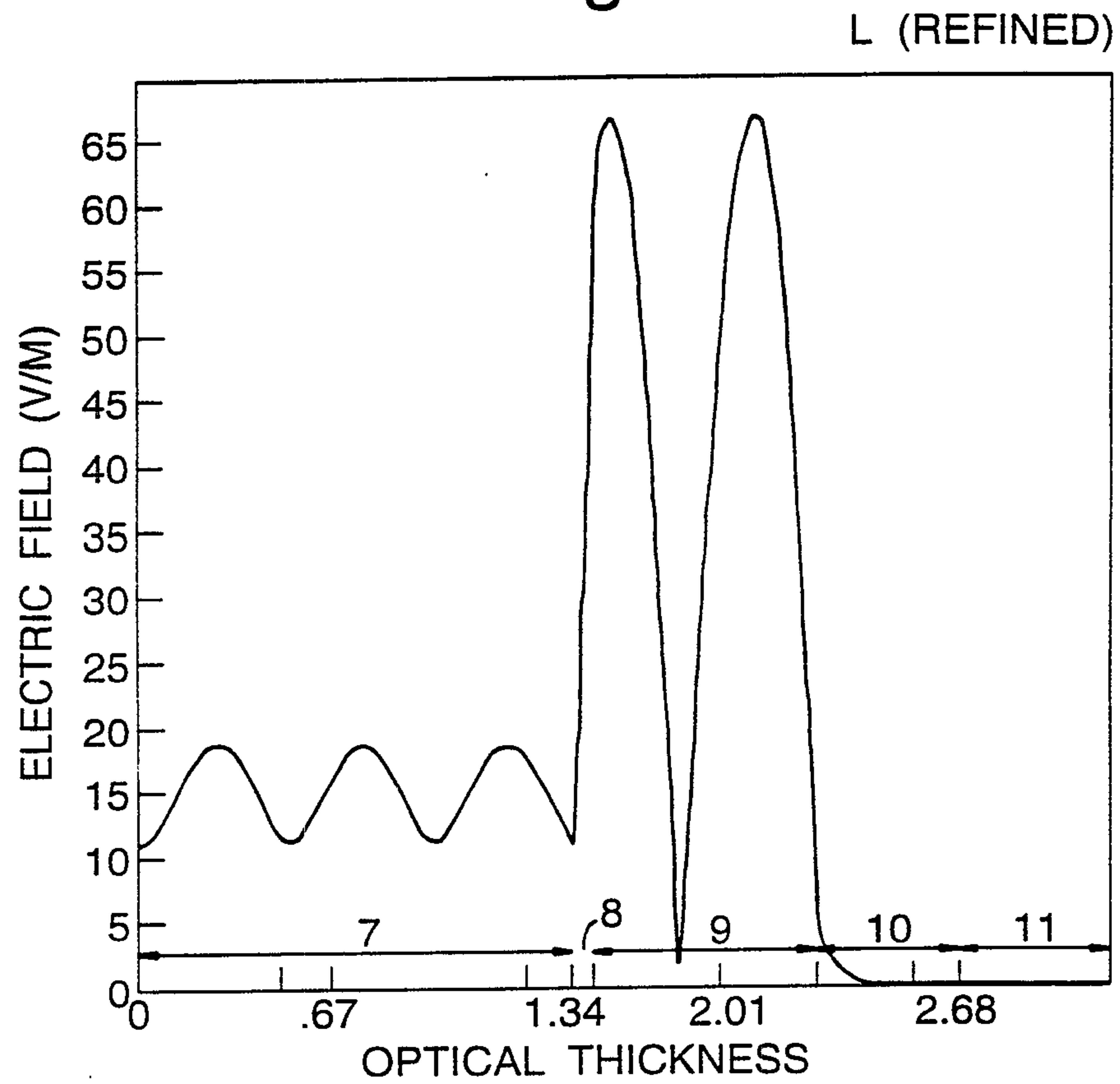
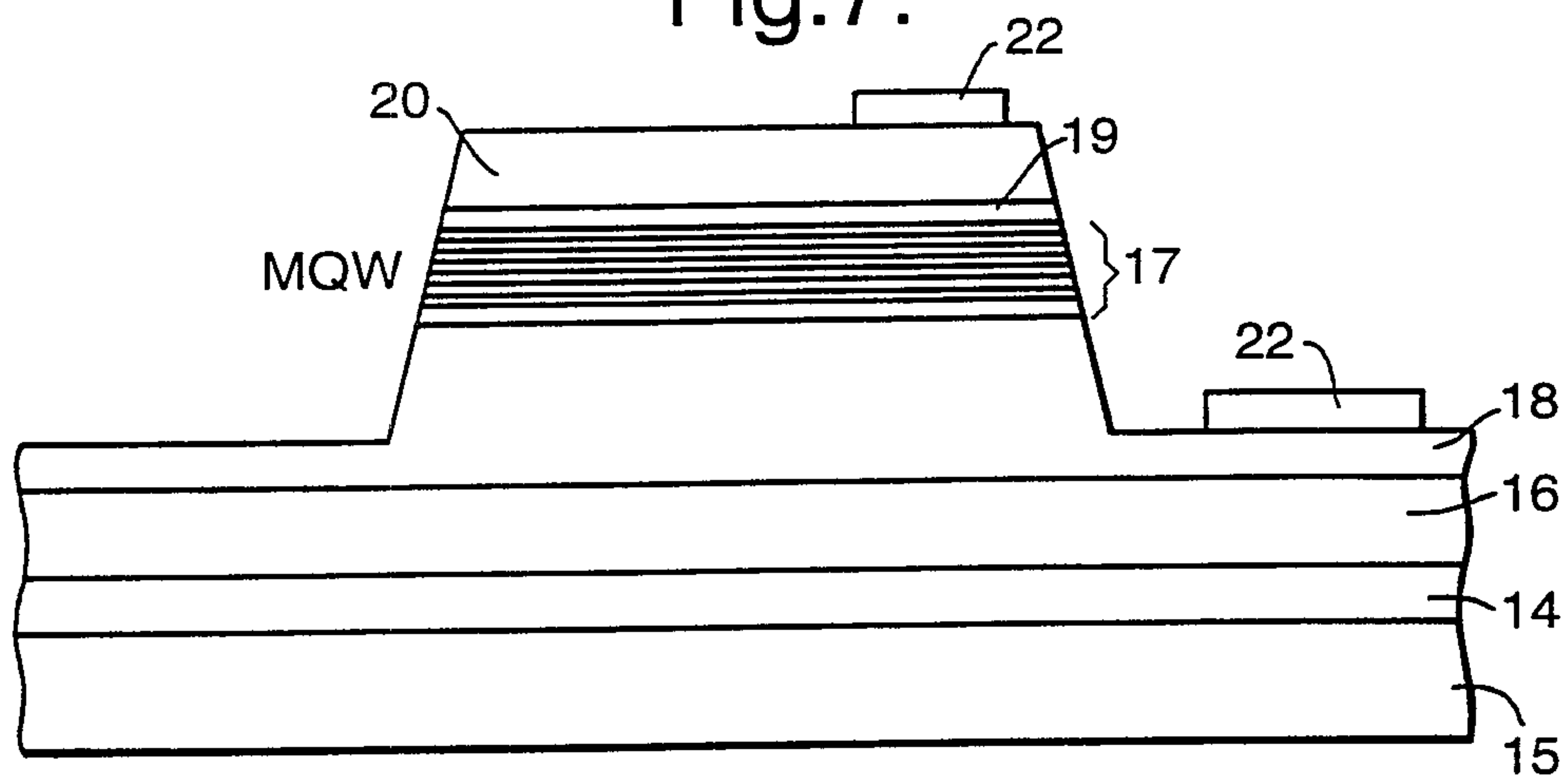
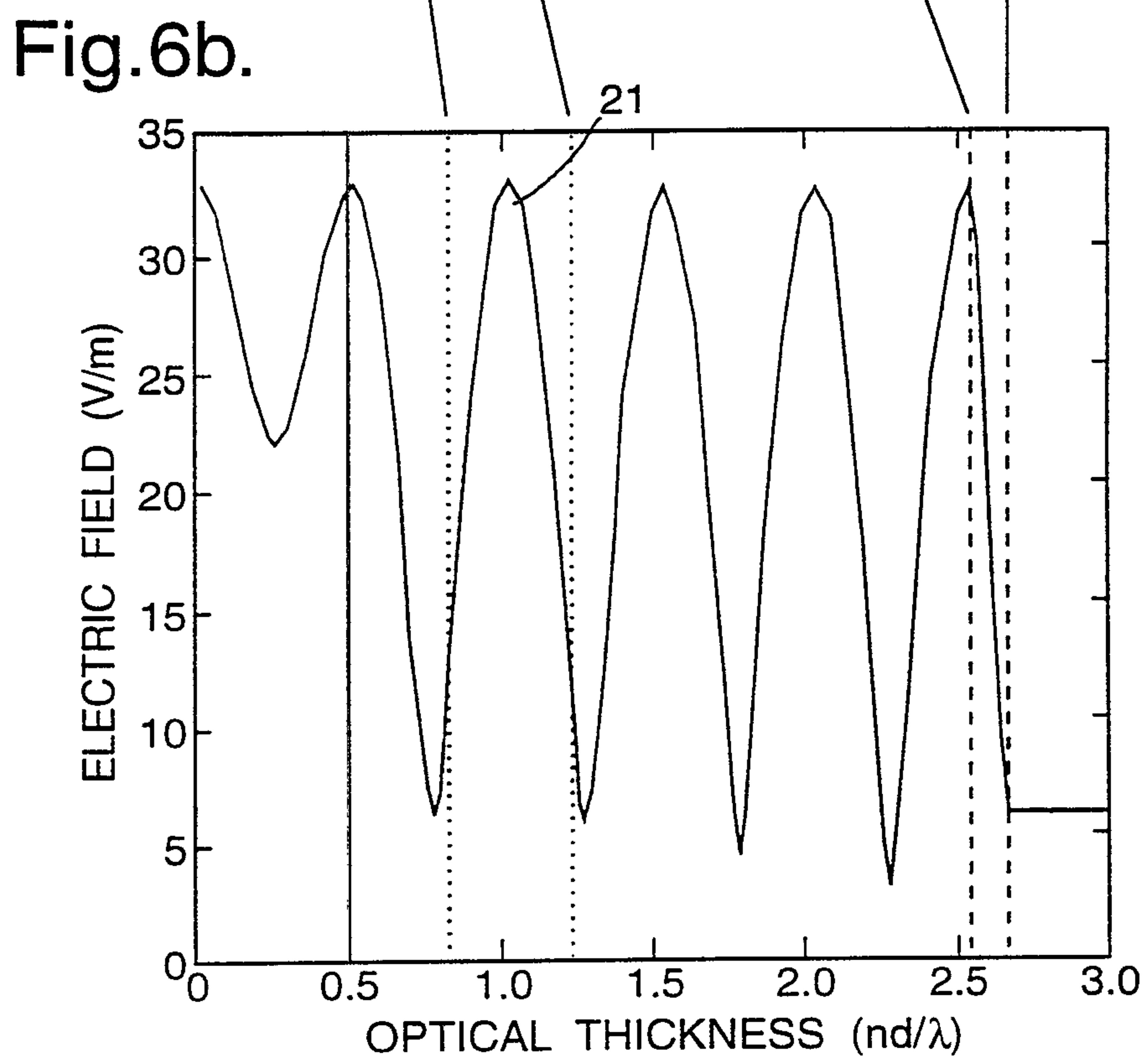
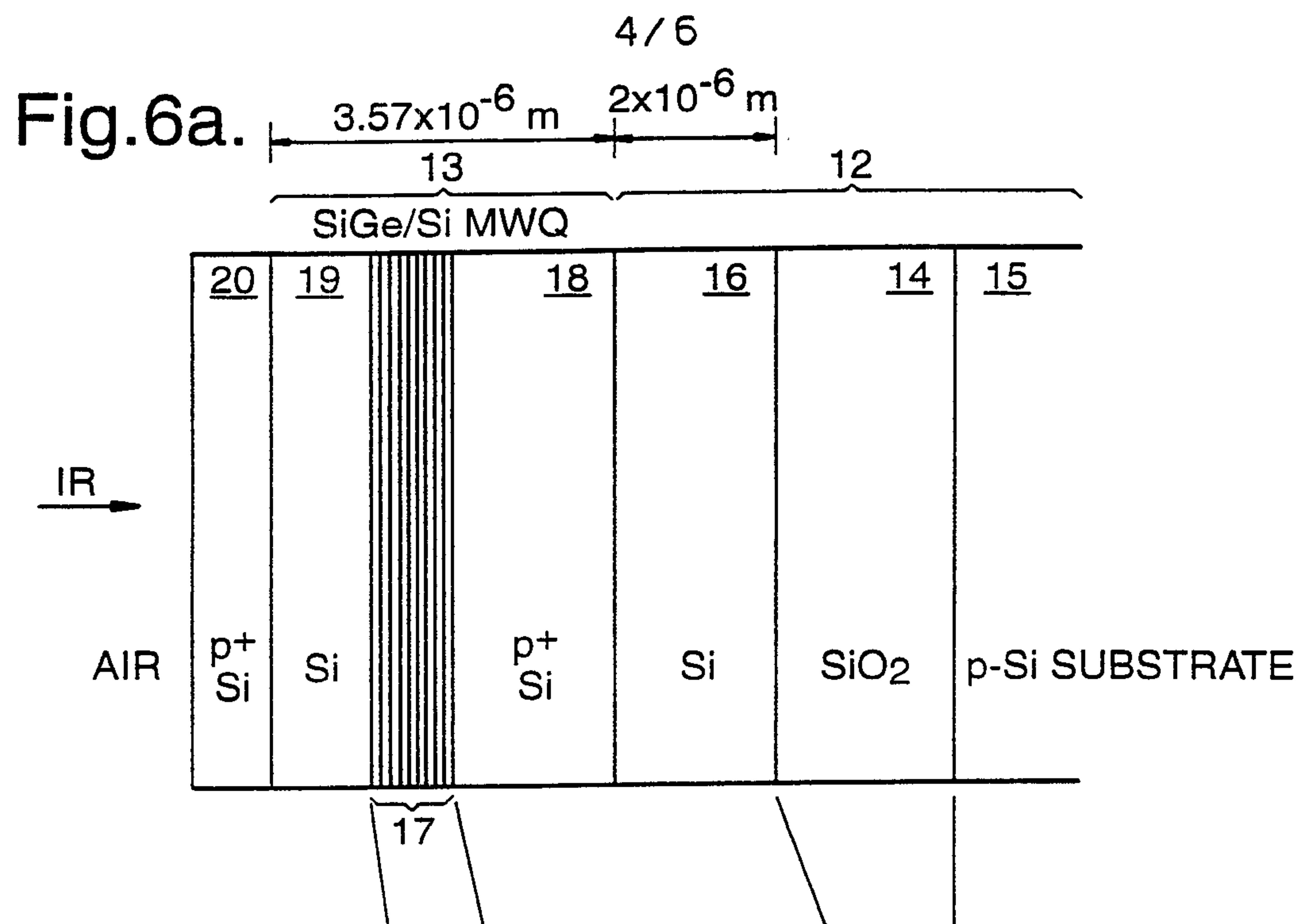


Fig.7.





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Fig.8.

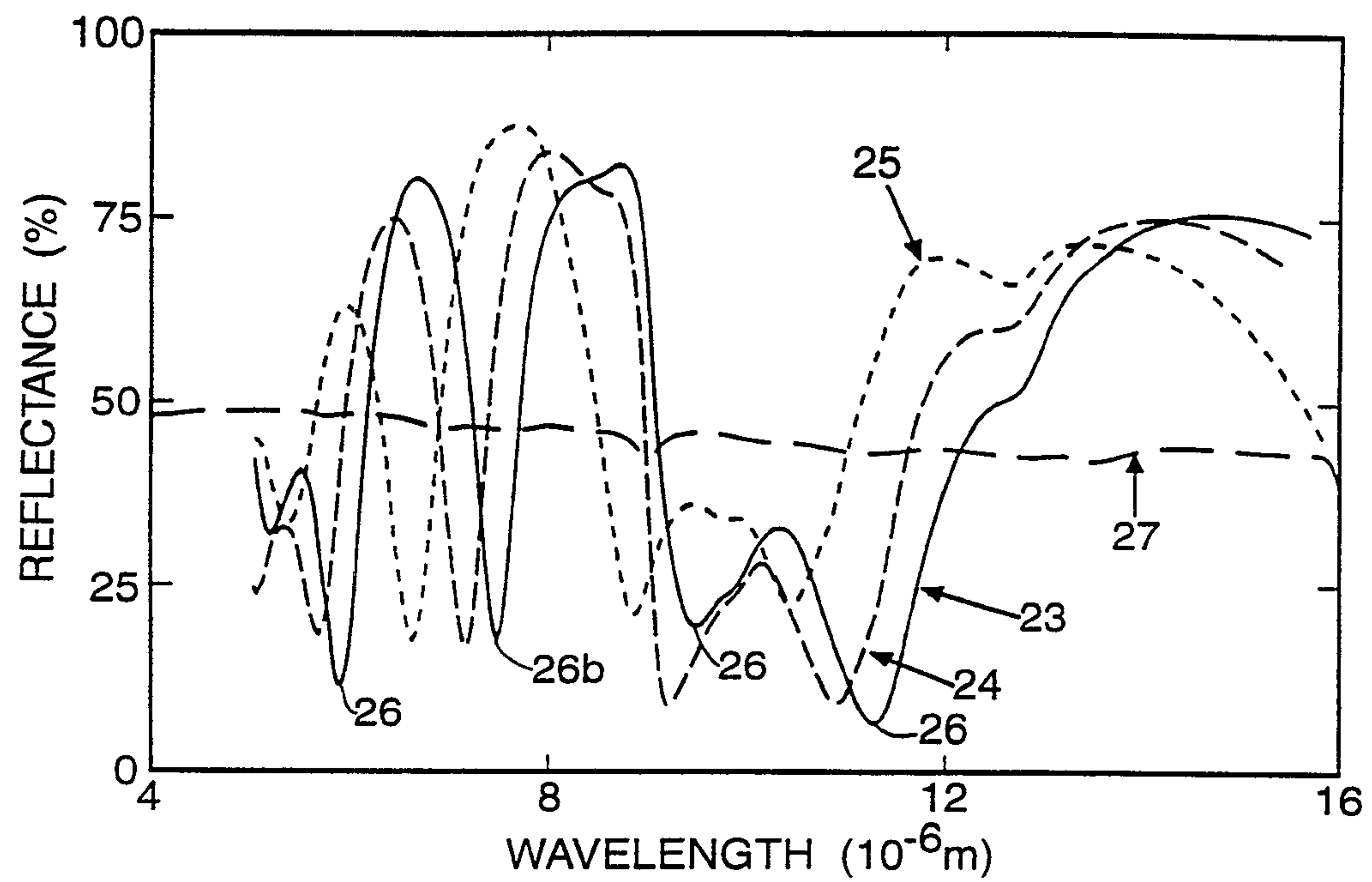
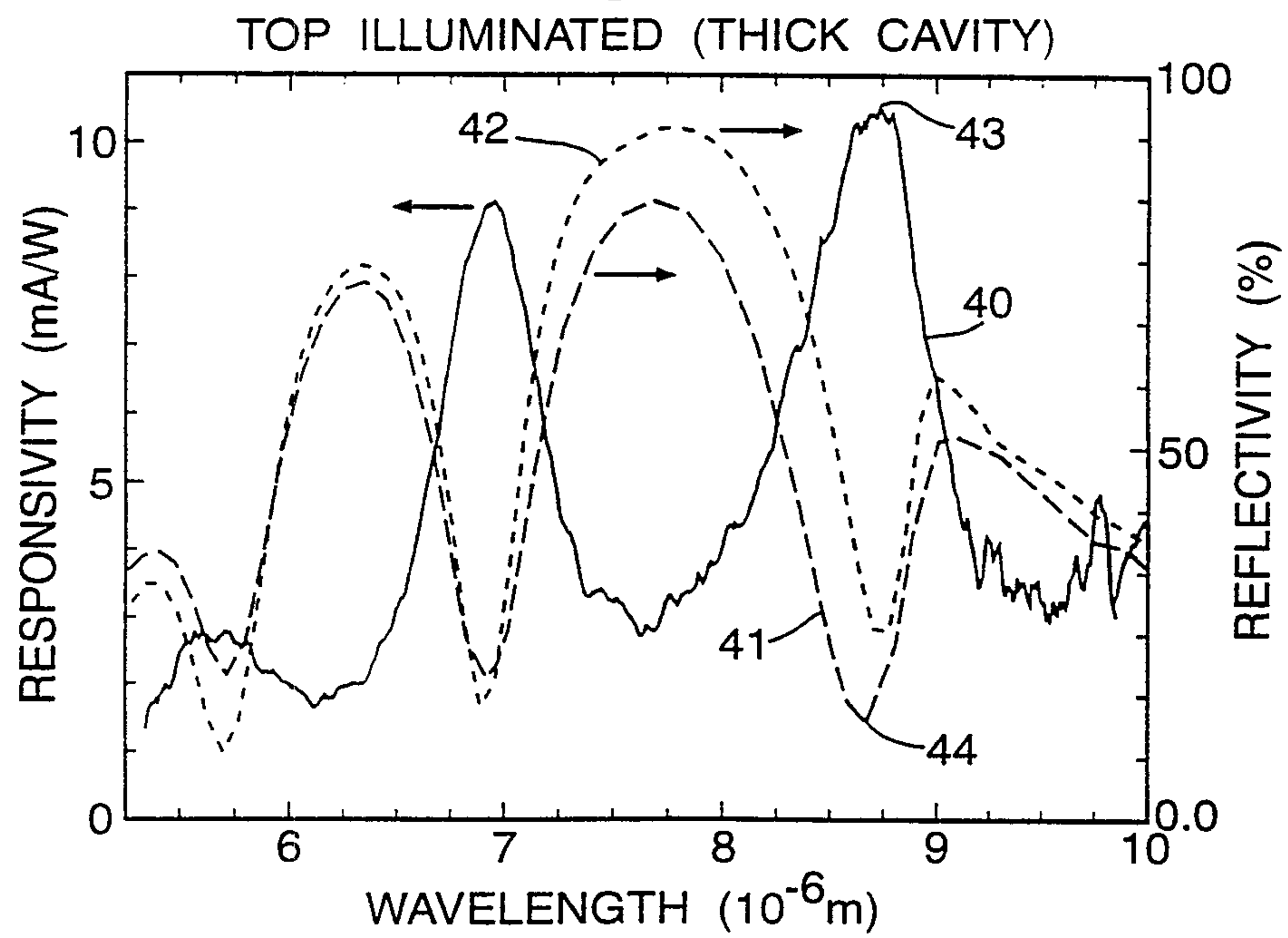


Fig.10.



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Fig.9a.

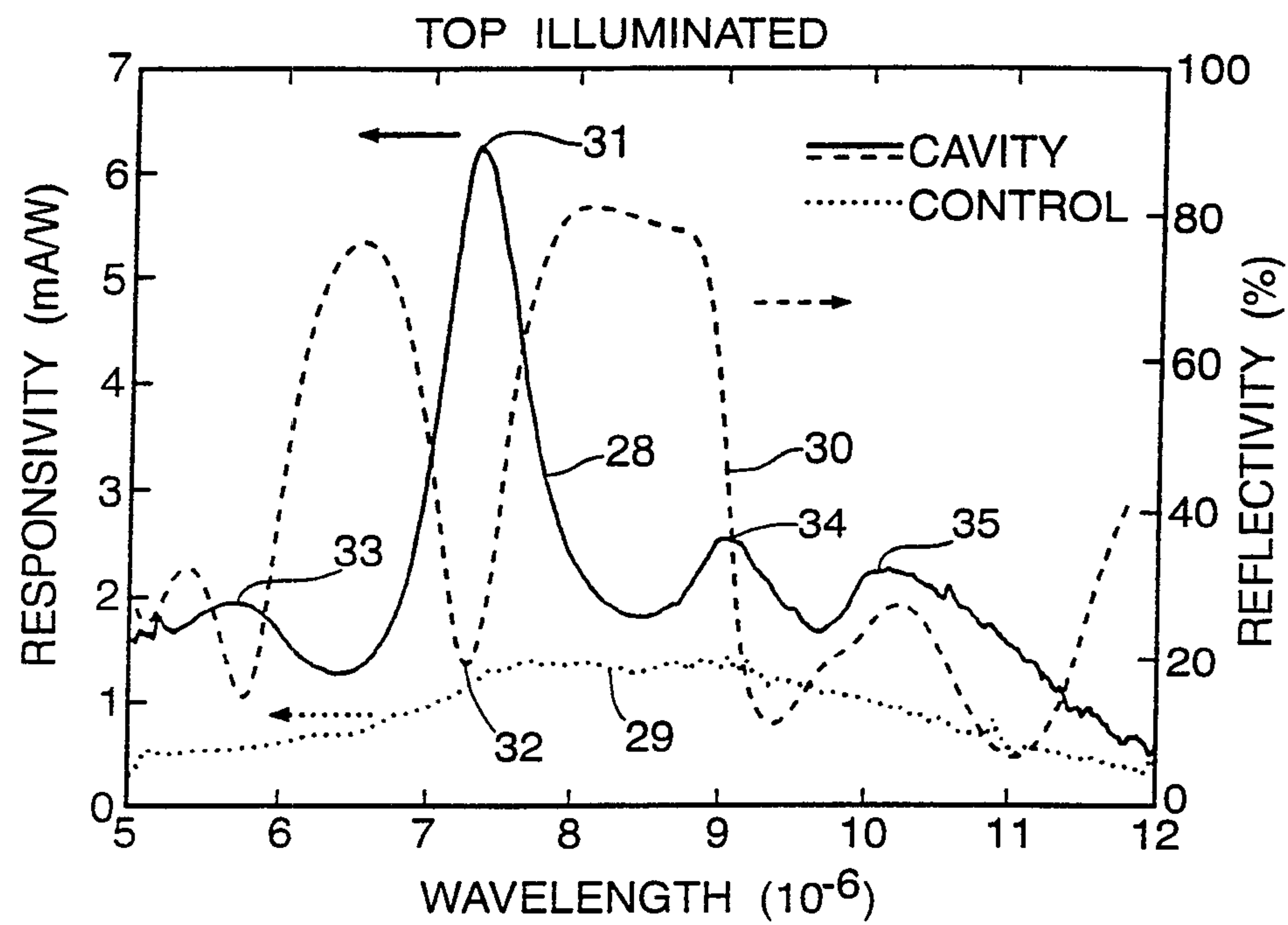


Fig.9b.

